Pressure-induced structural phase transitions and phonon anomalies in ReO₃: Raman and first-principles study

D. V. S. Muthu,¹ Pallavi Teredesai,^{1,2} S. Saha,^{1,3} Suchitra,⁴ U. V. Waghmare,⁴ A. K. Sood,^{1,5} and C. N. R. Rao⁵

¹Department of Physics, Indian Institute of Science, Bangalore 560012, India

²Navrachana University, Vadodara 391410, India

³Department of Physics, National University of Singapore, Singapore 117411

⁴Theoretical Science Unit, Jawaharlal Nehru Center for Advanced Scientific Research, Jakkur P.O., Bangalore 560064, India

⁵Chemistry and Physics of Materials Unit, New Chemistry Unit and International Centre for Materials Science, Jawaharlal Nehru Centre for

Advanced Scientific Research, Jakkur P.O., Bangalore 560064, India

(Received 17 February 2015; revised manuscript received 28 April 2015; published 19 June 2015)

We report high-pressure Raman-scattering studies on single-crystal ReO₃ up to 26.9 GPa at room temperature, complemented by first-principles density functional calculations to assign the modes and to develop understanding of the subtle features of the low-pressure phase transition. The pressure (*P*) dependence of phonon frequencies (ω) reveals three phase transitions at 0.6, 3, and 12.5 GPa with characteristic splitting and changes in the slope of $\omega(P)$. Our first-principles theoretical analysis confirms the role of the rotational modes of ReO₆, *M*₃, to the lowest pressure structural transition, and shows that the transition from the *Pm3m* to the *Im3* structure is a weak first-order transition, originating from the strong anharmonic coupling of the *M*₃ modes with the acoustic modes (strain).

DOI: 10.1103/PhysRevB.91.224308

PACS number(s): 78.30.-j, 63.20.dk, 64.60.-i

I. INTRODUCTION

 ABO_3 perovskite oxides exhibit a diverse range of phenomena ranging from magnetoresistance, superconductivity to ferroelectricity [1]. Typically, the insulating perovskite oxides, such as BaTiO₃, are known to commonly exhibit temperature-(T-) induced structural phase transitions in which the vanishing occupation of d orbitals of a transition-metal ion at the B site is a key to its off-centering as well as to the symmetry lowering structural instability. In contrast, T-dependent structural transitions are not common in metallic perovskite oxides, which typically have the d orbitals of their transition metal partially occupied by electrons. Competing structural instabilities in a perovskite involve rotations of the BO_6 octahedra, which are not excluded by the partial occupancy of d orbitals of the B cation [2]. When the size of the *B* cation relative to that of the *A* cation is large, the cubic perovskite structure can be unstable with respect to such octahedral rotational modes, and this can be another factor that can lead to structural transitions [1] and be used to understand fundamentally interesting structural transitions in metallic perovskite oxides. From this angle, ReO₃ is a rather interesting metallic perovskite oxide having conductivity within a factor of 6 of that of copper at room temperature and $\sim 100\%$ reflectivity below the plasma edge at 2.1 eV. Here, the A cation is completely missing, and four pressure-dependent (but interestingly no temperature-dependent) structural transitions are observed.

At ambient pressure, the structure of ReO_3 is cubic (space group Pm3m) consisting of corner-linked ReO_6 octahedra and linear Re-O-Re bonds. The empty A cation site in ReO_3 promotes structural instabilities involving rotations of the rigid ReO_6 octahedra. The early interest in ReO_3 was triggered by an observation of a supposedly continuous phase transition at ~ 0.5 GPa at room temperature with the compressibility being an order of magnitude higher in the high-pressure phase [3] (and hence termed as "compressibility collapse transition"). Early neutron-diffraction measurements showed that the high-pressure phase is tetragonal (P4/mbm), which transforms at a slightly higher pressure of 0.72 GPa to the cubic II (Im3) phase [4]. However, recent neutron-diffraction studies do not find any evidence of the tetragonal (P4/mbm)intermediate phase between the two cubic structures [5]. The softening of the M_3 phonon mode was suggested to be the driving force behind the Pm3m to Im3 pressure transition. X-ray diffraction at high pressures by Jorgensen et al. [6] showed three other phase transitions; a cubic II (Im3) phase to a monoclinic MnF₃ related phase at 3 GPa to the rhombohedral VF₃ related phase (rhombohedral I) at 12 GPa and to another rhombohedral phase (rhombohedral II) at 38 GPa. Subsequent high-pressure x-ray diffraction studies by Suzuki et al. [7] did not observe the monoclinic phase and showed that the cubic phase and rhombohedral I phase coexist in the pressure range of 8-18 GPa. Biswas et al. [8] have recently reported pressure-induced phase transitions in nanocrystalline ReO₃ using synchrotron x-ray diffraction. They have shown that the ambient pressure cubic I phase (Pm3m) changes to the monoclinic phase (C2/c) at ~0.3 GPa and from the monoclinic to a rhombohedral I phase (R3c) at ~6.7 GPa and finally to a rhombohedral II phase at ~20.3 GPa.

Purans *et al.* [9] have carried out Raman measurements at room temperature and ambient pressure. We carry out a high-pressure Raman study of crystalline ReO₃ here. We have carried out high-pressure Raman measurements up to 26.9 GPa on single crystals of ReO₃ showing three phase transitions at 0.6, 3, and 12.5 GPa. The pressure derivatives of the phonon frequencies have been determined. Our results for the high-pressure transitions corroborate the x-ray diffraction results of Jorgensen *et al.* [4]. However, the low-pressure structural transition at $P \sim 0.6$ GPa is quite intriguing, and we

^{*}Corresponding author: asood@physics.iisc.ernet.in